

# Investigation of Electromechanical Properties and Related Temperature Characteristics in Domain-Engineered Tetragonal Pb(In<sub>1/2</sub>Nb<sub>1/2</sub>) O<sub>3</sub>-Pb(Mg<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>-PbTiO<sub>3</sub> Crystals

Fei Li,<sup>‡,§</sup> Shujun Zhang,<sup>†,§</sup> Zhuo Xu,<sup>‡</sup> Xiaoyong Wei,<sup>‡</sup> Jun Luo,<sup>¶</sup> and Thomas R. Shrout<sup>§</sup>

<sup>‡</sup>Electronic Materials Research Laboratory, Key Laboratory of the Ministry of Education, Xi'an Jiaotong University, Xi'an 710049, China

<sup>§</sup>Materials Research Institute, Pennsylvania State University, University Park, Pennsylvania 16802

<sup>¶</sup>TRS Technologies Inc., State College, Pennsylvania 16801

The orientation-dependent electromechanical properties were calculated for tetragonal Pb(In<sub>1/2</sub>Nb<sub>1/2</sub>)O<sub>3</sub>-Pb(Mg<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>-PbTiO<sub>3</sub> (PIN-PMN-PT) crystals based on single-domain data. The maximum electromechanical coupling  $k_{33}^*$  was found to lie along the polar direction [001], whereas the maximum piezoelectric coefficient  $d_{33}^*$  was found to occur along [011]. Subsequently, the piezoelectric properties of [011] poled tetragonal PIN-PMN-PT crystals, with an engineered domain configuration ("2T"), were studied using resonance impedance measurement and strain versus electric field (S-E) behavior, where the piezoelectric coefficient  $d_{33}$  and coupling  $k_{33}$  of [011] poled crystals were found to be on the order of 1000 pC/N and 0.75, respectively. The high  $d_{33}$  of [011] poled crystals was associated with the high shear coefficient  $d_{15}$  (~2300 pC/N) in single domain state. Finally, the piezoelectric and electromechanical properties of [011] domain engineered tetragonal PIN-PMN-PT crystals were investigated as a function of temperature. In contrast to [001] single-domain PIN-PMN-PT crystals, the piezoelectric coefficient  $d_{33}$  and coupling  $k_{33}$ of [011] poled crystals were found to decrease with increasing temperature.

## I. Introduction

**R**<sub>Nb<sub>2/3</sub>)O<sub>3</sub>–PbTiO<sub>3</sub> (PMN–PT) and Pb(Zn<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>–PbTiO<sub>3</sub> (PZN–PT), have been extensively studied because of their ultra-high piezoelectric ( $d_{33} > 1500 \text{ pC/N}$ ) and electrome-chanical properties ( $k_{33} > 0.9$ ).<sup>1,2</sup> The ultra-high piezoelectric response in relaxor-based rhombohedral crystals with a "4R" engineered domain configuration, has been attributed to the polarization rotation reported by Fu and Cohen,<sup>3</sup> using first principle calculations. The usage temperature range of relaxor-PT-based crystals, however, is limited by their relatively low ferroelectric–ferroelectric phase transition  $T_{R-TS}$  (<100°C), which occurs at significantly lower temperatures than their respective Curie temperatures, due to strongly curved morphotropic phase boundarys (MPB).<sup>4</sup> Thus, extensive studies have been carried out in an attempt to increase  $T_c/T_{R-T}$ . The ternary crystal system Pb(In<sub>1/2</sub>Nb<sub>1/2</sub>)O<sub>3</sub>–Pb(Mg<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>–PbTiO<sub>3</sub> (PIN–PMN–PT) has been found to be a promising candidate, with compositions exhibiting  $T_{R-T}s > 120^{\circ}C$ , <sup>5–7</sup> 30°C higher than the phase-transition temperature of commercial binary PMN–PT crystals.</sub>

To further broaden the usage temperature range of relaxor-PT-based crystals, another approach is to utilize crystals with the tetragonal phase, where no ferroelectric–ferroelectric phase transition occurs before  $T_c$  and above room temperature. For the case of PIN–PMN–PT crystals,  $T_c$  is on the order of 220°C. Recently, a single-domain state was readily realized in [001] poled tetragonal PIN–PMN–PT crystals, with a relatively high electromechanical coupling  $k_{33}$ , being 0.84, maintaining this value to  $T_c$ .<sup>8</sup> Following the concept of domain engineering in rhombohedral crystals to achieve higher piezoelectric properties, it was proposed to investigate domain engineering in tetragonal PIN–PMN–PT crystals.

Based on the full set of material constants for single-domain tetragonal PIN–PMN–PT crystals,<sup>8</sup> the orientation dependence of the dielectric permittivity  $\varepsilon_{33}^*/\varepsilon_0$ , piezoelectric coefficient  $d_{33}^*$ , and electromechanical coupling  $k_{33}^*$  were calculated for tetragonal PIN–PMN–PT crystals at room temperature. The piezoelectric properties of [011] poled tetragonal crystals, with an engineered domain configuration ("2T"), were studied as a function of electric field and temperature and compared with single domain values.

# **II. Experimental Procedure**

Tetragonal PIN-PMN-PT single crystals were grown by the modified Bridgman technique. Details of the growth and properties of tetragonal PIN-PMN-PT crystals have been described in Li *et al.*<sup>8</sup>. The orientation dependence of the dielectric per-mittivity  $\varepsilon_{33}^*/\varepsilon_0$ , piezoelectric coefficient  $d_{33}^*$ , and electromechan-ical coupling factor  $k_{33}^*$  were calculated using the coordinate transformation method described in Nye and colleagues9,10 The PIN-PMN-PT crystals were oriented along crystallographic [001] and [011] directions by real-time Laue X-ray. The samples were electroded using sputtered gold thin film. The samples were poled using a dc field of 10 kV/cm at 120°C for 5 min, and subsequently field cooled to room temperature to avoid cracking. The resonance and antiresonance frequencies of longitudinal vibration mode for [011] and [001] poled tetragonal PIN-PMN-PT crystals were determined using an HP4194A impedance analyzer (Palo Alto, CA). The electromechanical coupling factors  $(k_{33})$  and piezoelectric coefficients  $(d_{33})$  were calculated following the IEEE standard.<sup>11</sup> The electric-field-induced strain measurements were carried out using a linear variable differential transducer (LVDT) driven by a lock-in amplifier (Mode SR830, Stanford Research System, Sunnyale, CA). The temperature dependence of the dielectric permittivity was determined from the capacitance using a multifrequency LCR meter (HP4284A), which was connected to a computer-controlled high-temperature furnace. The temperature dependence of the

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Table I. Piezoelectric and Dielectric Anisotropy of Single-Domain PMN-PT, PZN-PT, and PIN-PMN-PT Crystals

Relaxor crystals	Data for single-domain state							Engineered domain configuration
	$\epsilon_{33}/\epsilon_0$	$\epsilon_{11}/\epsilon_0$	$\epsilon_{11}/\epsilon_{33}$	<i>d</i> <sub>33</sub> (pC/N)	<i>d</i> <sub>15</sub> (pC/N)	$d_{15}/d_{33}$	k <sub>33</sub> (%)	k*33
PIMNT (T) <sup>†</sup>	1090	15000	14	530	2350	4	84	75% along [011]
PIMNT $(R)^{\ddagger}$	700	6300	9	74	2190	30	36	89% along [001]
PMN-42PT <sup>§</sup>	660	8600	13	260	1300	5	78	72% along [011]
PMN-33PT <sup>¶</sup>	640	3900	6	190	4100	21	69	94% along [001]
PZNT−12PT <sup>  </sup>	870	$12000^{\dagger}$	14	560	$1450^{\dagger}$	2.5	87	62% along [011]
PZNT-8PT <sup>††</sup>	1000	16000	16	$90^{\dagger}$	5000	55	39	94% along [001]

<sup>†</sup>This work. <sup>‡</sup>Liu et al.<sup>14</sup> <sup>§</sup>Cao et al.<sup>15</sup> <sup>¶</sup>Zhang et al.<sup>10</sup> <sup></sup><sup>¶</sup>Zhang et al.<sup>16</sup> <sup>††</sup>Park and Shrout<sup>1</sup> and Zhang et al.<sup>17</sup>

electromechanical coupling factors and piezoelectric coefficients were obtained using an impedance analyzer (HP4194A) connected to a temperature chamber.

## III. Results and Discussion

# (1) Anisotropic Behavior of Piezoelectric, Dielectric, and Electromechanical Properties

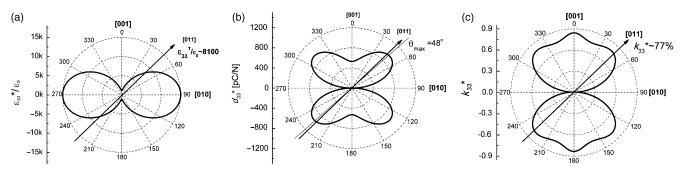
In order to determine the crystal orientations with optimal dielectric, piezoelectric, and electromechanical properties, it was deemed appropriate to study the level of anisotropic properties of tetragonal PIN–PMN–PT crystals. Table I gives the anisotropic properties for single-domain relaxor-PT-based crystals with rhombohedral and tetragonal phases. As shown in Table I, large dielectric ( $\varepsilon_{11}/\varepsilon_{33}$ ) and piezoelectric anisotropy factor ( $d_{15}/d_{33}$ ) were observed for tetragonal PIN–PMN–PT, being 14 and 4, respectively. These large anisotropic factors were also observed for other single-domain relaxor-PT-based crystals with compositions close to their respective MPBs. As reported by Budimir *et al.*,<sup>12,13</sup> large dielectric and piezoelectric anisotropic factors are due to the proximity of the ferroelectric– ferroelectric phase transition, whether induced by changes in composition or temperature, or by the application of an external electric field.

From the single-domain data of tetragonal PIN-PMN-PT crystals,<sup>8</sup> the orientation dependence of the dielectric permittivity  $\varepsilon_{33}^*/\varepsilon_0$ , piezoelectric coefficient  $d_{33}^*$ , and electromechanical coupling factor  $k_{33}^*$  were calculated and presented in Fig. 1. It was found that the dielectric permittivity  $\epsilon_{33}^*/\epsilon_0$  increased as the orientation deviated away from the [001] polar direction, due to the  $\varepsilon_{11}$  dielectric permittivity being larger than  $\varepsilon_{33}$  for tetragonal crystals. As shown in Fig. 1(b), the highest value of the piezoelectric coefficient  $d_{33}^*$  was on the order of 950 pC/N, being along the direction 48° away from the spontaneous polarization (close to [011] direction). According to Damjanovic and colleagues, 18,19 the largest longitudinal piezoelectric response will be found away from the polar direction when the ratio  $d_{15}/d_{33}$  is larger than 1.5 ( $Q_{1133}/Q_{3333}$ ), i.e., the critical value of piezoelectric anisotropy  $(d_{15}/d_{33} = 1.5)$ , where Q is the electrostrictive coefficient. Based on the critical value, crystals can be categorized into "rotator" and "extender" ferroelectrics.<sup>19</sup> For "extender" ferroelectrics, the value of  $d_{15}/d_{33}$  is smaller than the critical value and the piezoelectric response is dominated by the collinear piezoelectric effect due to the extension of the polar vector; thus, the longitudinal coefficient  $d_{33}^*$  is highest along the polar direction. For "rotator" ferroelectrics, however, the value of  $d_{15}/d_{33}$  is higher than the critical value, with polarization rotation associated with the shear piezoelectric effect, being the dominant mechanism contributing to the piezoelectric response. As a result,  $d_{33}^*$  reaches its peak value along a direction away from the spontaneous polarization.<sup>19</sup> Based on the concept above, the tetragonal PIN–PMN–PT crystals belong to "rotator" ferroelectrics at room temperature.

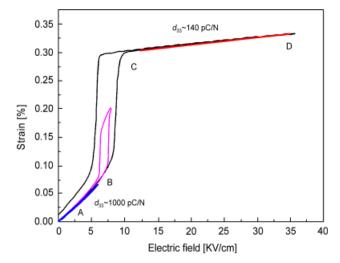
In contrast to  $d_{33}^*$ , the maximum coupling factor  $k_{33}^*$  was found to lie along the [001] polar direction, as shown in Fig. 1(c), demonstrating that domain engineering would not be beneficial to electromechanical couplings. This phenomenon is different from that observed in rhombohedral crystals, where domain engineering benefits both piezoelectric coefficients and coupling factors.<sup>1,10</sup> To delineate this difference, the piezoelectric and dielectric anisotropy for single-domain tetragonal and rhombohedral crystals were compared. As summarized in Table I, tetragonal crystals exhibit a much larger anisotropic behavior in the dielectric property when compared with the piezoelectric property, while rhombohedral crystals exhibit a larger piezoelectric anisotropy. This tendency is also observed for other oxygen-octahedra ferroelectric systems,19 being attributed to the higher electrostrictive ratio  $2Q_{1313}/Q_{3333}$  in rhombohedral crystals when compared with tetragonal ones. For example, for the case of the nonpolar [011] direction in tetragonal PIN-PMN–PT crystals, the dielectric permittivity  $\varepsilon_{33}^*/\varepsilon_0$  was found to be eight times higher than the value along the [001] polar direction, whereas the piezoelectric coefficient  $d_{33}^*$  increased only by 80%. Consequently, the electromechanical coupling  $k_{33}^*$  along the [011] direction is calculated to be only 0.77  $(k_{33}^3 = d_{33}^*/\sqrt{\epsilon_{33}^* s_{33}^*})$ , where  $s_{33}^*$  were 23 and 41 pm<sup>2</sup>/N along the [011] and [001] directions, respectively), which is smaller than the value along the [001] direction, as shown in Fig. 1(c).

# (2) Piezoelectric Response of [011] Poled Tetragonal PIN–PMN–PT Crystals

As demonstrated above, the maximum piezoelectric coefficient  $d_{33}^*$  and coupling factor  $k_{33}^*$  were found to be along the crystallographic [011] and [001] directions, respectively. In this section,



**Fig.1.** Orientation dependence of (a)  $\varepsilon_{13}^*/\varepsilon_0$ , (b)  $d_{33}^*$ , and (c)  $k_{13}^*$  for single domain tetragonal PIN–PMN–PT crystals in the [001]–[010] plane.



**Fig. 2.** Strain versus electric field behavior for PIN–PMN–PT single crystals poled along the [011] direction at various electric fields.

the piezoelectric properties for [011] domain-engineered tetragonal PIN–PMN–PT crystals were investigated and compared with the [001] poled single domain counterparts.

The piezoelectric coefficients  $d_{33}$  for [001] and [011] poled crystals, determined using the resonance method, were found to be 530 and 1050 pC/N, respectively. As shown in Fig 1(b), the calculated piezoelectric coefficient  $d_{33}^*$  along the [011] direction was found to be 940 pC/N, which is very close to the experimental result. Therefore, the measured high  $d_{33}$  for [011] poled crystal mainly arises from an intrinsic contribution, being associated with the high  $d_{15}$  of single domain property. The electromechanical coupling factors  $k_{33}$  were found to be 0.84 and 0.75 for [001] and [011] poled crystals, respectively, corresponding to the calculated results, as shown in Fig. 1(c), demonstrating that the contribution of domain wall is minimal.

The mechanical quality factor along [001] and [011] directions were calculated and found to be on the order of 700 and 250, respectively. Crystallographic-dependent mechanical quality factor was also reported for rhombohedral PMN–PT crystals, in which higher quality factors were observed along the spontaneous polarization [111] direction, owing to the existence of pesudomonodomain state.<sup>20</sup>

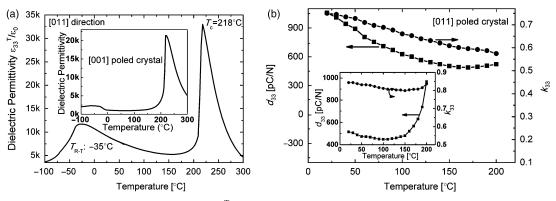
Figure 2 shows a strain–electric field (*S*–*E*) loop for [011] poled samples at various levels of electric fields. Tetragonal PIN– PMN–PT crystals poled along the [011] direction showed an engineered domain configuration (noted as "2T"); therefore, the *S*– *E* loop was hysteresis-free in the low field section (the portion AB in Fig. 2). The piezoelectric coefficient  $d_{33}$  calculated from the slope of *S*–*E* loop was found to be on the order of 1000 pC/N, corresponding to the results obtained by the resonance technique, again demonstrating that the domain wall contribution to the piezoelectric response was minimal. A discontinuous jump in the strain curve was observed when the electric field increased to 7 kV/cm, corresponding to a tetragonal—orthorhombic phase transition.<sup>21</sup> This electric-field-induced phase transition was thought to be first-order in nature, due to the large *S*–*E* hysteresis. At the higher field portion, "CD" in the *S*–*E* loop, the strain was saturated and linear, with a piezoelectric coefficient calculated to be 140 pC/N. This value corresponds to the crystal being in an orthorhombic phase, where the applied electric field was along the polarization direction and the piezoelectric response was only contributed by the polarization extension.<sup>19</sup>

# (3) Temperature-Dependent Dielectric, Piezoelectric, and Electromechanical Properties for Single-Domain and Domain-Engineered Tetragonal Crystals

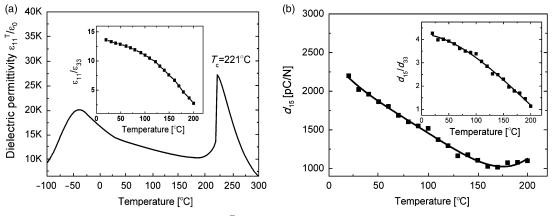
Figure 3(a) shows the temperature dependence of the dielectric permittivity  $\varepsilon_{33}^1/\varepsilon_0$  for [011] and [001] poled crystals. Unlike the temperature-independent behavior in [001] poled crystals, (shown in the inset of Fig. 3(a)), the dielectric permittivity  $\varepsilon_{33}^{T}/\varepsilon_{0}$ of [011] poled crystals was found to decrease with an increasing temperature from room temperature to 170°C, with the variation being -30%. Fig. 3(b) shows the temperature dependence of the piezoelectric coefficients and coupling factors for [011] and [001] poled crystals, where the piezoelectric coefficient  $d_{33}$  of [011] poled crystals was found to follow a similar trend. Owing to the large degradation of  $d_{33}$  (decreased by 50%), the electromechanical coupling  $k_{33}$  of [011] poled crystals was found to decrease by 25% at 170°C, when compared with the roomtemperature value. In contrast, the temperature-independent behavior of electromechanical coupling  $k_{33}$  and piezoelectric coefficient  $d_{33}$  was found in the temperature range of  $20^{\circ}$ – $170^{\circ}$ C for [001] poled crystals.

To better understand the degradation of dielectric and piezoelectric properties in [011] poled crystals (2T) with an increasing temperature, the temperature dependences of  $d_{15}$  and  $\varepsilon_{11}^{T}/\varepsilon_0$  for [001] poled single-domain crystals were measured. As given in Fig. 4, the dielectric permittivity  $\varepsilon_{11}^{T}/\varepsilon_0$  and piezoelectric coefficient  $d_{15}$  of [001] poled crystals were found to decrease by 30% and 55% respectively, in the temperature range of 20°–170°C. Therefore, the degradation of  $\varepsilon_{33}^{T}/\varepsilon_0$  and  $d_{33}$  for [011] poled crystals was mainly attributed to the reduction of  $\varepsilon_{11}^{T}/\varepsilon_0$  and  $d_{15}$  in single-domain state. As reported for tetragonal BaTiO<sub>3</sub> crystals, the degradations of  $\varepsilon_{11}^{T}/\varepsilon_0$  and  $d_{15}$  were attributed to the test temperature being away from the ferroelectric–ferroelectric phase transition temperature<sup>12</sup> ( $T_{\text{FF}}$  is about –35°C for the investigated tetragonal PIN–PMN–PT crystals).

The temperature-dependent dielectric and piezoelectric anisotropy factors for tetragonal PIN–PMN–PT crystals were also determined, as shown in the inset of Figs. 4(a) and (b), respectively. The dielectric anisotropy factor  $(\epsilon_{11}^T/\epsilon_{33}^T)$  decreased from 14 to 3, while the piezoelectric anisotropy factor  $(d_{15}/d_{33})$  was found to decrease from 4.2 to 1, with the increasing temperature.



**Fig. 3.** Temperature dependence of (a) dielectric permittivity  $\varepsilon_{33}^{T}/\varepsilon_0$  (measured at 1 kHz), (b) piezoelectric coefficient  $d_{33}$  and coupling  $k_{33}$  for [011] and [001] (small insets) poled tetragonal PIN–PMN–PT crystals.



**Fig. 4.** Temperature dependence of (a) dielectric permittivity  $\varepsilon_{11}^{T}/\varepsilon_0$  (measured at 1 kHz) and (b) piezoelectric coefficient  $d_{15}$  for [001] poled single domain tetragonal PIN–PMN–PT crystals. The dielectric and piezoelectric anisotropy factors are presented in the insets of (a) and (b), respectively.

Consequently, tetragonal PIN–PMN–PT crystals became "extender" ferroelectrics at elevated temperatures where domain engineering cannot benefit the piezoelectric response.

# IV. Conclusion

The orientation dependence of dielectric, piezoelectric, and electromechanical properties was calculated based on the single-domain data. The maximum electromechanical coupling  $k_{33}^*$  was found to be along the spontaneous polarization direction [001], whereas the maximum piezoelectric coefficient  $d_{33}^*$  was found to lie along the [011] direction. In contrast to domain-engineered rhombohedral crystals, domain engineering in tetragonal PIN-PMN-PT crystals did not lead to enhanced electromechanical coupling factors, owing to their large dielectric anisotropy ( $\varepsilon_{11}$ /  $\varepsilon_{33}$ ). The piezoelectric and dielectric properties in domainengineered tetragonal PIN-PMN-PT crystals were found to decrease with increasing temperature, being associated with the dielectric permittivity  $\varepsilon_{11}^{T}/\varepsilon_{0}$  and piezoelectric coefficient  $d_{15}$  in single-domain state being degraded as the temperature moved away from ferroelectric-ferroelectric transition temperature.

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